



Phonon-polariton mediated heat transfer in nanostructures

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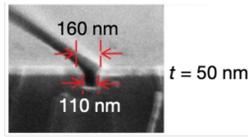
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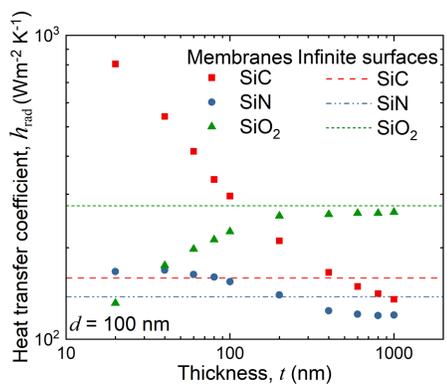
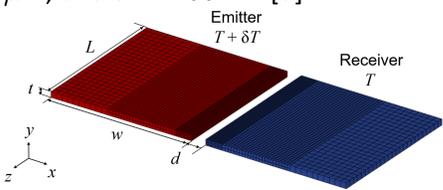


Radiative heat transfer in the dual nanoscale regime

Different polaritonic materials with similar EM responses display opposite trends in the dual nanoscale regime of near-field radiative heat transfer (NFRHT) [1, 2].



Using numerically-exact discrete system Green's function (DSGF) simulations [3, 4], the NFRHT coefficient is predicted between membranes with $w = L = 1 \mu\text{m}$, and $d = 100 \text{ nm}$ [5]:



$$h_{\text{rad}}(T) = \int_0^\infty \frac{1}{2\pi A_c} \left[\sum_{i \in V_E} \sum_{j \in V_R} \mathcal{T}_{ij}(\omega) \right] \left[\frac{\partial \Theta(\omega, T')}{\partial T} \right]_{T'=T} d\omega,$$

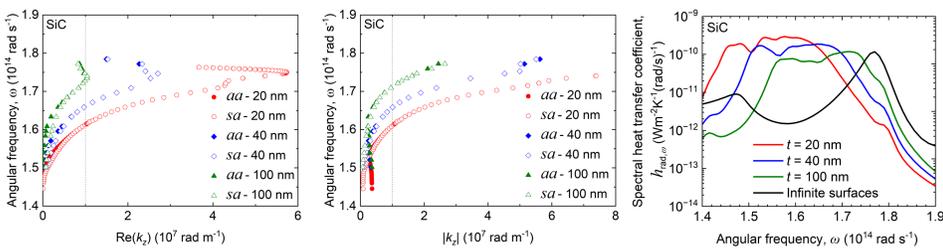
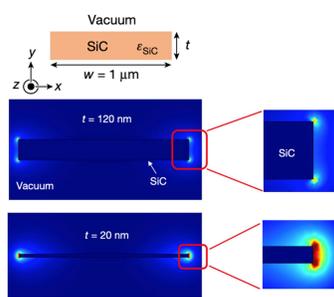
where Θ is the mean energy of an EM state, and \mathcal{T}_{ij} is the transmission coefficient.

Modal analyses show that all membranes support **EM corner and edge modes** [6], which induce spectral redshift.

Material losses (i.e., $\text{Im}(\epsilon)$) play a key role in the dual nanoscale regime.

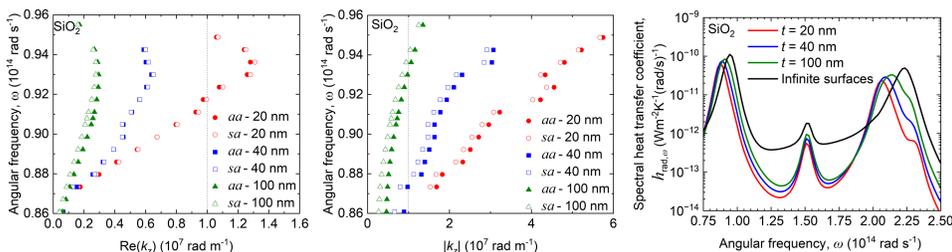
SiC:

- Losses are mostly negligible, $\text{Re}(k_z) \gg \text{Im}(k_z)$.
- Either $\text{Re}(k_z)$ or $|k_z|$ can predict $h_{\text{rad},\omega}$ resonances.
- Strong SPhP coupling is promoted, **enhancing** NFRHT.



SiN and SiO₂:

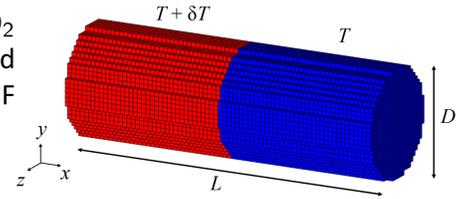
- Losses are non-negligible.
- $|k_z|$ should be used to predict resonance redshift.
- Losses reduce the maximum $\text{Re}(k_z)$ tunneled between the membranes.
- For SiO₂, NFRHT is **attenuated** with decreasing membrane thicknesses.



Application: localized radiative cooling.

Radiation-conduction in nanowires

The radiative conductivity in SiO₂ nanowires with $D = 66 \text{ nm}$ and $L = 200 \text{ nm}$ is predicted using the DSGF method [3, 7]:



$$\kappa(T) = \int_0^\infty \frac{1}{2\pi A_c} \left[\sum_{i \in V_A} \sum_{j \in V_B} \mathcal{T}_{ij}(\omega) (\mathbf{r}_j - \mathbf{r}_i) \cdot \hat{\mathbf{x}} \right] \left[\frac{\partial \Theta(\omega, T')}{\partial T} \right]_{T'=T} d\omega.$$

The radiative conductivity is typically calculated using kinetic theory (KT) as [8]:

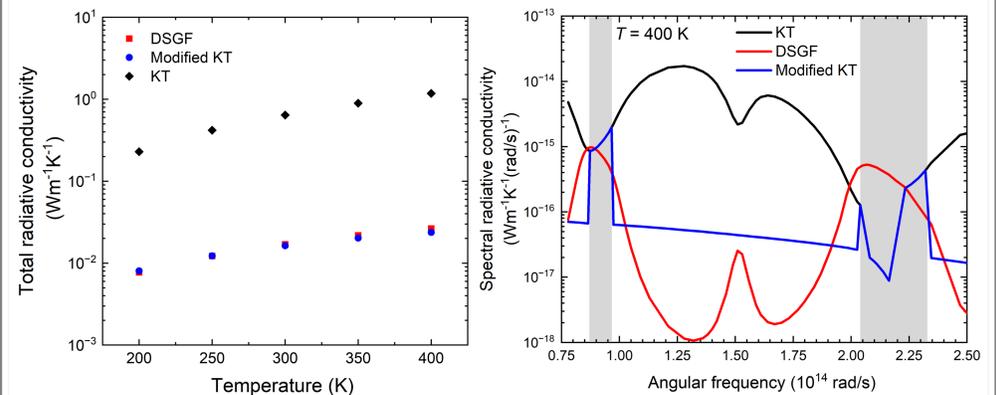
$$\kappa(T) = \int_0^\infty \frac{1}{\pi A_c} \hbar \omega \Lambda \frac{\partial f_0}{\partial T} d\omega,$$

where the mean free path Λ is calculated using dispersion relations for an infinite cylinder.

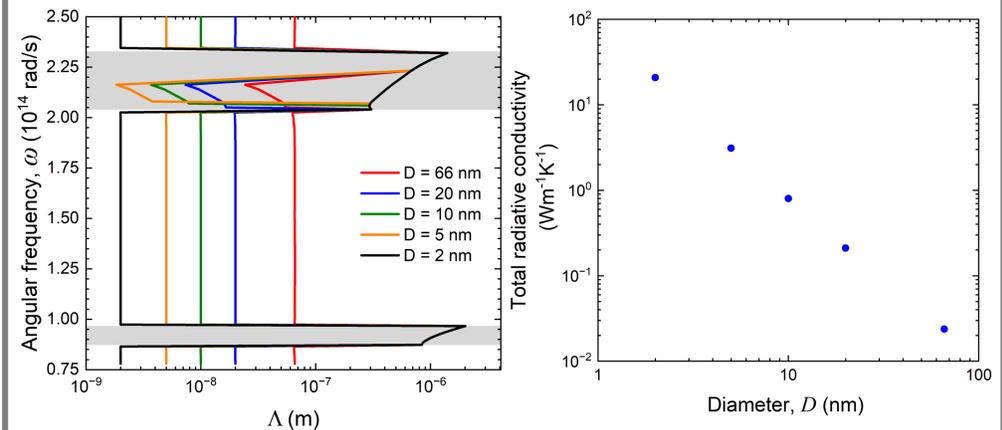
Assuming Λ as the propagation length $L_p = (2\text{Im}(k))^{-1}$ is not valid over the whole spectrum.

We propose a modified KT where

- Inside Reststrahlen bands: $\Lambda = L_p$. SPhPs are dominant and do not undergo size effect.
- Outside Reststrahlen bands: $\Lambda = \Lambda_{\text{wire}} \left(\frac{1}{\text{Kn}} - \frac{1}{4\text{Kn}^2} \right)$, where $\Lambda_{\text{wire}} = L_p$ and $\text{Kn} = \Lambda_{\text{wire}}/D$. Photons experience **size effect** due to boundary scattering.



Radiation-conduction can exceed the bulk phonon conductivity when SiO₂ nanowires are thinner than 10 nm.



Application: thermal management of solid-state devices.

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